# Computational Study of the Release of $\mathbf{H}_{2}$ from Ammonia Borane Dimer $\left(\mathrm{BH}_{3} \mathrm{NH}_{3}\right)_{2}$ and Its Ion Pair Isomers 

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Received: April 25, 2007; In Final Form: June 4, 2007


#### Abstract

High-level electronic structure calculations have been used to map out the relevant portions of the potential energy surfaces for the release of $\mathrm{H}_{2}$ from dimers of ammonia borane, $\mathrm{BH}_{3} \mathrm{NH}_{3}(\mathbf{A B})$. Using the correlationconsistent aug-cc-pVTZ basis set at the second-order perturbation MP2 level, geometries of stationary points were optimized. Relative energies were computed at these points using coupled-cluster $\operatorname{CCSD}(\mathrm{T})$ theory with the correlation-consistent basis sets at least up to the aug-cc-pVTZ level and in some cases extrapolated to the complete basis set limit. The results show that there are a number of possible dimers involving different types of hydrogen-bonded interactions. The most stable gaseous phase $(\mathbf{A B})_{2}$ dimer results from a head-totail cyclic conformation and is stabilized by $14.0 \mathrm{kcal} / \mathrm{mol}$ with respect to two $\mathbf{A B}$ monomers. ( $\mathbf{A B})_{2}$ can generate one or two $\mathrm{H}_{2}$ molecules via several direct pathways with energy barriers ranging from 44 to 50 $\mathrm{kcal} / \mathrm{mol}$. The diammoniate of diborane ion pair isomer, $\left[\mathrm{BH}_{4}^{-}\right]\left[\mathrm{NH}_{3} \mathrm{BH}_{2} \mathrm{NH}_{3}{ }^{+}\right]$(DADB), is $10.6 \mathrm{kcal} / \mathrm{mol}$ less stable than $(\mathbf{A B})_{2}$ and can be formed from two $\mathbf{A B}$ monomers by overcoming an energy barrier of $\sim 26$ $\mathrm{kcal} / \mathrm{mol}$. DADB can also be generated from successive additions of two $\mathrm{NH}_{3}$ molecules to $\mathrm{B}_{2} \mathrm{H}_{6}$ and from condensation of $\mathbf{A B}$ with separated $\mathrm{BH}_{3}$ and $\mathrm{NH}_{3}$ molecules. The pathway for $\mathrm{H}_{2}$ elimination from DADB is characterized by a smaller energy barrier of $20.1 \mathrm{kcal} / \mathrm{mol}$. The alternative ion pair $\left[\mathrm{NH}_{4}{ }^{+}\right]\left[\mathrm{BH}_{3} \mathrm{NH}_{2} \mathrm{BH}_{3}{ }^{-}\right]$is calculated to be $16.4 \mathrm{kcal} / \mathrm{mol}$ above $(\mathbf{A B})_{2}$ and undergoes $\mathrm{H}_{2}$ release with an energy barrier of $17.7 \mathrm{kcal} /$ mol. $\mathrm{H}_{2}$ elimination from both ion pair isomers yields the chain $\mathrm{BH}_{3} \mathrm{NH}_{2} \mathrm{BH}_{2} \mathrm{NH}_{3}$ as product. Our results suggest that the neutral dimer will play a minor role in the release of $\mathrm{H}_{2}$ from ammonia borane, with a dominant role from the ion pairs as observed experimentally in ionic liquids and the solid state.


## Introduction

There is considerable interest in the development and application of hydrogen-based fuel cells as a viable alternative energy source based on renewable fuels. A critical issue with hydrogen as a fuel for use in on-board transportation systems is the need to find efficient chemical $\mathrm{H}_{2}$ storage materials that enable the safe and efficient release/uptake of $\mathrm{H}_{2} .{ }^{1}$ Ammonia borane $\left(\mathrm{BH}_{3} \mathrm{NH}_{3}, \mathbf{A B}\right)^{2-4}$ and its derivatives have recently emerged as promising hydrogen sources. ${ }^{5-11}$ Using high accuracy ab initio electronic structure theory calculations of reaction thermodynamic parameters, in both the gas phase and the solid state, we have shown that molecular $\mathbf{A B}$ and the corresponding salt $\left[\mathrm{BH}_{4}^{-}\right]\left[\mathrm{NH}_{4}^{+}\right]$can serve as good hydrogen storage systems that release $\mathrm{H}_{2}$ in slightly exothermic processes. ${ }^{6}$ The isovalent systems including alane amine $\left(\mathrm{AlH}_{3} \mathrm{NH}_{3}\right)$, borane phosphine $\left(\mathrm{BH}_{3} \mathrm{PH}_{3}\right)$, and phosphine alane $\left(\mathrm{AlH}_{3} \mathrm{PH}_{3}\right)$ and their corresponding salts $\left[\mathrm{XH}_{4}^{-}\right]\left[\mathrm{YH}_{4}{ }^{+}\right]$, with $\mathrm{X}=\mathrm{B}$ or Al and $\mathrm{Y}=$ N or P , have also been shown to possess appropriate thermodynamic properties to serve as $\mathrm{H}_{2}$ storage systems. ${ }^{7}$

The potential of a compound for chemical hydrogen storage is dependent on the inherent kinetics and mechanism of the processes releasing $\mathrm{H}_{2}$ and regenerating the compound that stores $\mathrm{H}_{2}$. Using thermo-analytical techniques, Wolf and co-

[^0]workers ${ }^{5}$ found that $\mathbf{A B}$ begins to undergo stepwise thermal decompositions at temperatures as low as 410 K . During the first decomposition step, approximately 1 mol of $\mathrm{H}_{2}$ per mole of $\mathbf{A B}$ is released, together with aminoborane $\left(\mathrm{BH}_{2} \mathrm{NH}_{2}\right)$, and other polymeric derivatives. Subsequent quantum chemical calculations ${ }^{12}$ using various composite methods suggested that formation of the first $\mathrm{H}_{2}$ from an $\mathbf{A B}$ molecule occurs through a concerted mechanism characterized by a planar four-membered transition structure with an energy barrier of around $32-33 \mathrm{kcal} /$ mol. In a recent much higher-level computational study, ${ }^{13}$ we showed that, in its monomeric form, $\mathbf{A B}$ undergoes preferentially $\mathrm{B}-\mathrm{N}$ bond cleavage giving $\mathrm{BH}_{3}$ and $\mathrm{NH}_{3}$. In fact, at the CCSD(T)/complete basis set (CBS) level of theory, the $\mathrm{B}-\mathrm{N}$ bond dissociation energy of $26 \mathrm{kcal} / \mathrm{mol}$ is smaller than the energy barrier of $37 \mathrm{kcal} / \mathrm{mol}$ for the concerted $\mathrm{H}_{2}$ elimination. With a difference of $11 \mathrm{kcal} / \mathrm{mol}$ in barrier heights, the thermal decomposition of $\mathbf{A B}$ is dominated by the bond cleavage channel. Under these conditions, $\mathrm{H}_{2}$ formation is not competitive, inconsistent with the experimental results mentioned above. We showed that a borane molecule could effectively act as a bifunctional Lewis acid catalyst favoring $\mathrm{H}_{2}$ formation from $\mathbf{A B}$. With $\mathrm{BH}_{3}$ formed by breaking the $\mathrm{B}-\mathrm{N}$ bond in $\mathbf{A B}$, the corresponding energy barrier is thus reduced to about $6 \mathrm{kcal} /$ mol with respect to the $\mathrm{BH}_{3} \mathrm{NH}_{3}+\mathrm{BH}_{3}$ separated reactants. ${ }^{13}$ These results show that without a catalyst the rate of $\mathrm{H}_{2}$ release from the $\mathbf{A B}$ monomer is very low at low temperatures. We note that the starting $\mathbf{A B}$ materials exist as a solid.

TABLE 1: Relative Energies ( $\mathbf{k c a l} / \mathrm{mol}$ ) of Ammonia Borane Dimer and Its Ion Pair Isomers Calculated at the $\operatorname{CCSD}(\mathrm{T})$ Level with Different Basis Sets ${ }^{a}$

| structure | aVDZ | aVTZ | aVQZ | CBS |
| :--- | :---: | ---: | :---: | ---: |
| dim | 0.0 | 0.0 | 0.0 | 0.0 |
| DADB | 10.6 | 10.7 | 10.6 | 10.6 |
| IonP-N | 16.5 | 16.6 | 16.5 | 16.4 |

${ }^{a}$ Using the aug-cc-pVnZ basis sets, with $n=\mathrm{D}, \mathrm{T}$, and Q , and extrapolated CBS energies. Relative energies include zero-point corrections obtained from MP2/aVDZ harmonic vibrational frequencies scaled by a factor of 0.9751 .

TABLE 2: Relative Energies (kcal/mol) of Stationary Points Related to $\mathbf{H}_{\mathbf{2}}$ Release from Ammonia Borane Dimers dim and DADB Calculated at the CCSD(T) Level with Different Basis Sets ${ }^{a}$

| structure | aVDZ | aVTZ | aVQZ | CBS |
| :--- | ---: | ---: | ---: | ---: |
| 2AB | 0.0 | 0.0 | 0.0 | 0.0 |
| dim | -14.2 | -14.0 | -13.7 | -13.5 |
| dim-ts1 | 36.3 | 36.0 | 35.9 | 35.9 |
| dim-p1 $+2 \mathrm{H}_{2}$ | -13.8 | -15.4 | -16.0 | -16.4 |
| 2( $\left.\mathrm{BH}_{2} \mathrm{NH}_{2}\right)+2 \mathrm{H}_{2}$ | -10.7 | -12.6 | -13.5 | -14.0 |
| DADB | -3.5 | -3.3 | -3.0 | -2.8 |
| DADB-ts1 | 17.2 | 17.2 | 17.2 | 17.3 |
| DADB-p1 $+\mathrm{H}_{2}$ | -7.8 | -9.3 | -9.4 | -9.3 |

${ }^{a}$ On the basis of $\operatorname{CCSD}(\mathrm{T})$ energies using MP2/aVTZ-optimized geometries and ZPE corrections from scaled MP2/aVDZ harmonic vibrational frequencies.

It has experimentally been found that "seeding" of $\mathbf{A B}$ results in significantly faster $\mathrm{H}_{2}$ production. The seed compound has been proposed to be the diammoniate of diborane $\left[\mathrm{BH}_{4}{ }^{-}\right]$ $\left[\mathrm{NH}_{3} \mathrm{BH}_{2} \mathrm{NH}_{3}{ }^{+}\right]$(DADB), an ion pair isomer of the ( $\left.\mathbf{A B}\right)_{2}$ dimer. ${ }^{14}$ Recent in situ solid-state ${ }^{11} \mathrm{~B}$ magic-angle spinning NMR studies ${ }^{15}$ at $88^{\circ} \mathrm{C}$ demonstrated that during the thermal decomposition of $\mathbf{A B}$ to produce $\mathrm{H}_{2}$ DADB is present. These authors suggested that the formation of DADB is the prerequisite for $\mathrm{H}_{2}$ release. In ionic liquids containing $\mathbf{A B}$, the formation of DADB was also observed. ${ }^{8 b}$ These authors also suggested an important role for DADB in $\mathrm{H}_{2}$ release noting that there was no induction period for $\mathrm{H}_{2}$ release in the ionic liquid whereas there was an induction period in the pure solid. The ionic liquid would favor the formation of ionic intermediates such as DADB. It has been reported that DADB produces polyaminoborane upon heating ${ }^{16}$ and that polyaminoboranes were observed in the ionic liquid as $\mathrm{H}_{2}$ was released. Another starting point connecting to the ion pair is that of diborane and two ammonia molecules $\left(\mathrm{B}_{2} \mathrm{H}_{6}+2 \mathrm{NH}_{3}\right) .{ }^{17}$ Another ion pair $\left[\mathrm{NH}_{4}{ }^{+}\right]\left[\mathrm{BH}_{3} \mathrm{NH}_{2} \mathrm{BH}_{3}{ }^{-}\right]$ (IonP-N) could also exist as an alternative isomeric form of $(\mathrm{AB})_{2}$.

In the present work, we have carried out high-level electronic structure calculations to map out the relevant portions of the potential energy surface of the $(\mathbf{A B})_{2}$ dimer. Our goal is to probe further the molecular mechanism of the $\mathrm{H}_{2}$ release from $\mathrm{BNH}_{6}$ materials. We report the existence of a range of reaction pathways from the $(\mathbf{A B})_{2}$ dimer leading to $\mathrm{H}_{2}$ production.

## Computational Methods

Calculations were performed using the Gaussian $03^{18}$ and MOLPRO $2002^{19}$ suites of programs. Geometry parameters of the stationary points were initially optimized using second-order perturbation theory (MP2) ${ }^{20}$ in conjunction with the correlationconsistent aug-cc-pVDZ basis set. ${ }^{21}$ The character of each stationary point was subsequently determined to be an equilib-


Figure 1. Selected MP2/aVTZ geometry parameters of the most stable ammonia borane dimer dim and its ion pair isomers DADB and IonP$\mathbf{N}$. Each structure is projected in two different plans. Bond lengths are given in $\AA$, and bond angles in degrees.

TABLE 3: Calculated Heats of Formation at 0 and 298 K (kcal/mol) for $\mathbf{B}-\mathbf{N}$ Compounds ${ }^{a}$

| molecule | $\Delta H_{\mathrm{f}}(0 \mathrm{~K})$ | $\Delta H_{\mathrm{f}}(298 \mathrm{~K})$ |
| :---: | :---: | :---: |
| $\mathrm{NH}_{3}{ }^{\text {b }}$ | $-9.6 \pm 0.5$ | $-11.3 \pm 0.5$ |
| $\mathrm{NH}_{4}{ }^{\text {b }}$ | $153.6 \pm 0.5$ | $150.9 \pm 0.7$ |
| $\mathrm{BH}_{3}{ }^{\text {b }}$ | $26.4 \pm 0.7$ | $25.5 \pm 0.7$ |
| $\mathrm{BH}_{4}{ }^{-b}$ | $-11.6 \pm 0.7$ | $-13.5 \pm 0.7$ |
| $\mathrm{BH}_{2} \mathrm{NH}_{2}{ }^{\text {b }}$ | $-15.9 \pm 1.0$ | $-18.6 \pm 1.0$ |
| $\mathrm{BH}_{3} \mathrm{NH}_{3}{ }^{\text {b }}$ | $-9.1 \pm 1.0$ | $-13.5 \pm 1.0$ |
| $\mathrm{BH}_{3} \mathrm{NH}_{2} \mathrm{BH}_{3}{ }^{-c}$ | $-49.4 \pm 1.2$ | $-55.3 \pm 1.2$ |
| $\mathrm{NH}_{3} \mathrm{BH}_{2} \mathrm{NH}_{3}{ }^{+c}$ | $105.9 \pm 1.2$ | $99.2 \pm 1.2$ |
| $2\left(\mathrm{BH}_{3} \mathrm{NH}_{3}\right)(\mathrm{dim})^{d}$ | $-32.2 \pm 1.0$ | $-41.3 \pm 1.0$ |
| $\mathrm{BH}_{3} \mathrm{NH}_{2} \mathrm{BH}_{2} \mathrm{NH}_{3}$ (DADB-p1) ${ }^{\text {d }}$ | $-27.4 \pm 1.0$ | $-35.3 \pm 1.0$ |
| $\mathrm{BH}_{3} \mathrm{NH}_{2} \mathrm{BH}_{2} \mathrm{NH}_{3}\left(\right.$ DADB-p3) ${ }^{\text {d }}$ | $-39.2 \pm 1.0$ | $-47.6 \pm 1.0$ |
| $\mathrm{NH}_{3} \mathrm{BH}_{2} \mathrm{NH}_{3} \cdots \mathrm{BH}_{4}(\mathbf{D A D B})^{c}$ | $-21.6 \pm 1.0$ | $-31.0 \pm 1.0$ |
| $\mathrm{BH}_{3} \mathrm{NH}_{2} \mathrm{BH}_{3} \cdots \mathrm{NH}_{4}(\mathbf{I o n P - N})^{c}$ | $-15.8 \pm 1.0$ | $-25.4 \pm 1.0$ |

${ }^{a}$ From $\operatorname{CCSD}(\mathrm{T}) / \mathrm{CBS}+$ correction calculations. ${ }^{b}$ Reference 6. ${ }^{c}$ This work. ${ }^{d}$ Reference 42.
rium structure or a first-order saddle point, by second derivative calculations yielding harmonic vibrational frequencies at the same level. Harmonic vibrational frequencies were scaled to estimate the zero-point energies (ZPEs) based on previous analyses. ${ }^{6,7,13}$ We used the accurate ZPE of $\mathrm{BH}_{3} \mathrm{NH}_{3}{ }^{6}$ as a reference value to determine the scaling factor for $\left(\mathrm{B}_{x} \mathrm{~N}_{x} \mathrm{H}_{y}\right)$

dim-ts1
Figure 2. Selected MP2/aVTZ geometry parameters of the transition state structure dim-ts1 for release of two $\mathrm{H}_{2}$ molecules from dim. Bond lengths are given in $\AA$, and bond angles in degrees.
derivatives. For MP2/aVDZ harmonic frequencies, a scaling factor of 0.9751 was obtained as the average of the MP2/aVDZ plus the accurate ZPE divided by the MP2/aVDZ value for $\mathrm{BH}_{3}-$ $\mathrm{NH}_{3}$. To ascertain the identity of the relevant transition state (TS) structures, intrinsic reaction coordinate (IRC) ${ }^{22}$ calculations were carried out using density functional theory with the hybrid B3LYP functional ${ }^{23}$ and the $6-311++G(d, p)$ basis set. The IRC calculations were done with the default step size of $0.1 \mathrm{amu}^{1 / 2}$ bohr were calculated. Geometry parameters of the relevant structures were then refined at the MP2/aug-cc-pVTZ level. ${ }^{21}$

Improved relative energies between stationary points were calculated using coupled-cluster theory at the $\operatorname{CCSD}(\mathrm{T})^{24}$ level using the MP2/aug-cc-pVTZ-optimized geometries. As a calibration for the variation of relative energies with respect to the basis set expansions, the $\operatorname{CCSD}(\mathrm{T})$ total energies of a few important structures were first extrapolated to the CBS limit, based on electronic energies obtained with the correlationconsistent aug-cc-pVnZ basis sets, with $n=\mathrm{D}, \mathrm{T}$, and Q . The extrapolation of the $\operatorname{CCSD}(\mathrm{T})$ energies to the CBS limit was done using the following expression ${ }^{25}$
$E(n)=A_{\mathrm{CBS}}+B \exp [-(n-1)]+C \exp \left[-(n-1)^{2}\right]$
with $n=2,3$, and 4 for the aug-cc-pVnZ, $n=\mathrm{D}, \mathrm{T}$, and Q , basis sets, respectively (referred to hereafter as aVnZ). For the remaining structures, only $\operatorname{CCSD}(\mathrm{T}) / \mathrm{aVTZ}$ energies were considered. We have previously found that use of the aVTZ results will change the relative energies by less than $1 \mathrm{kcal} /$ mol. ${ }^{13,26}$ In the MP2 and $\operatorname{CCSD}(\mathrm{T})$ calculations, the core orbitals were kept frozen.

The lattice energy of the ionic salts was estimated using the volume approach proposed by Jenkins et al. ${ }^{27}$ following previous work from the Bartlett group. ${ }^{28}$ It is calculated from the empirical expression

$$
\begin{equation*}
U_{\mathrm{L}}=2 I\left[\alpha V_{\mathrm{m}}^{-1 / 3}+\beta\right] \tag{2}
\end{equation*}
$$

where $I$ is the ionic strength ( $I=1$ for MX (1:1) salts), and $V_{\mathrm{m}}$ is the molecular (formula unit) volume of the lattice involved, which is equal to the sum of the individual ion volumes of the cation, $V_{+}$, and anion, $V_{-}$. The electron densities were calculated using the B3LYP/DZVP2 level of density functional theory, and the volume corresponds to that inside the 0.001 au contour of the electron density. The choice of contour level was made on the basis of volumes used in free energy of solvation calculations. ${ }^{29}$ For the empirical parameters in eq 2, the following values were used: $\alpha=28.0 \mathrm{kcal} / \mathrm{mol}$ and $\beta=12.4$ $\mathrm{kcal} / \mathrm{mol} .{ }^{27}$ The lattice energies can be corrected to standard conditions at 298 K using eq $3^{27}$


Figure 3. Schematic potential energy profiles, in $\mathrm{kcal} / \mathrm{mol}$, showing the reaction pathways for $\mathrm{H}_{2}$ release from two ammonia borane monomers $2\left(\mathrm{BH}_{3} \mathrm{NH}_{3}\right)$ via the dimer $\operatorname{dim}$ and from DADB, calculated with the $\operatorname{CCSD}(\mathrm{T})$ method using four basis sets including ZPE corrections. The energies are from top to bottom: aVDZ, aVTZ, aVQZ, and CBS.


Figure 4. Selected MP2/aVTZ geometry parameters of four different TS's for release of one $\mathrm{H}_{2}$ molecule from dim. Bond lengths are given in $\AA$, and bond angles in degrees.

$$
\begin{equation*}
U_{\mathrm{L}}(298 \mathrm{~K})=U_{\mathrm{L}}+\left[p\left(n_{\mathrm{M}} / 2-2\right)+q\left(n_{\mathrm{X}} / 2\right)\right] R T \tag{3}
\end{equation*}
$$

for $\mathrm{M}_{p} \mathrm{X}_{q}$ salts, where $n_{\mathrm{M}}=n_{\mathrm{X}}=6$ for nonlinear polyatomic ions.

## Results and Discussion

Stability of Ammonia Borane Dimer and Its Ion Pair Isomers. The calculated relative energies of the three isomeric structures are listed in Table 1. The complexation energy of the dimer with respect to the two $\mathbf{A B}$ monomers, chosen as the energy reference, is summarized in Table 2. Total CCSD(T)


Figure 5. Schematic potential energy profiles illustrating five different reaction pathways for $\mathrm{H}_{2}$ release from the ammonia borane dimer dim. Relative energies given in $\mathrm{kcal} / \mathrm{mol}$ were obtained from $\operatorname{CCSD}(\mathrm{T}) / \mathrm{aVTZ}$ + ZPE calculations.
and ZPEs for the isomers and the $\mathrm{NH}_{3} \mathrm{BH}_{2} \mathrm{NH}_{3}{ }^{+}$and $\mathrm{BH}_{3} \mathrm{NH}_{2} \mathrm{BH}_{3}{ }^{-}$ ions are tabulated in Table S-1 of the Supporting Information.

The geometry and complexation energy of the $(\mathbf{A B})_{2}$ dimer has previously been investigated. ${ }^{30-42}$ The nature of the hydrogen-bond interaction in the dimer has been the subject of most of these analyses. Our present calculations confirm that the dimer has a number of low-energy conformations with different symmetry point groups. The energy differences between these different conformations are very small, less than $0.5 \mathrm{kcal} / \mathrm{mol}$. The structure characterized by $C_{2 h}$ symmetry and denoted hereafter as dim, is the energetically lowest-lying dimer conformation that we found. One $\mathrm{N}-\mathrm{H}$ bond of an $\mathbf{A B}$ monomer interacts with two $\mathrm{B}-\mathrm{H}$ bonds of the other monomer. Selected optimized geometry parameters at the MP2/aVTZ level are displayed in Figure 1. The $\mathrm{H}(\mathrm{N}) \cdots \mathrm{H}(\mathrm{B})$ intermolecular distance of $1.986 \AA$ in dim, optimized at MP2/aVTZ level, differs from the values optimized at lower levels. ${ }^{30-34}$ The BH and $\mathrm{N}-\mathrm{H}$ bond distances within the $\mathrm{B}-\mathrm{H} \cdots \mathrm{H}-\mathrm{N}$ dihydrogen bond are longer than the non-interacting ones by up to $0.01 \AA$. The atomic net charges $(q)$ of the four $\mathrm{B}-\mathrm{H} \cdots \mathrm{H}-\mathrm{N}$ atoms are: $q(\mathrm{~B})=0.42 e, q(\mathrm{H})=-0.33 e, q(\mathrm{H})=0.20 e$, and $q(\mathrm{~N})=$ $0.17 e$ at the $\mathrm{HF} / \mathrm{aVTZ}$ level, corresponding to a $\mathrm{B}-\mathrm{H}^{\delta-} \cdots \mathrm{H}-\mathrm{N}^{\delta+}$ charge distribution consistent with the analysis of Li et al. ${ }^{37}$ These authors suggested in addition that the main factor of the hydrogen bond is a charge transfer from the occupied $\sigma(\mathrm{B}-\mathrm{H})$ orbital to the vacant $\sigma^{*}(\mathrm{X}-\mathrm{H})$ orbitals. ${ }^{37}$

We have previously calculated the heat of formation of dim $\Delta H_{\mathrm{f}}\left[\left(\mathrm{BH}_{3} \mathrm{NH}_{3}\right)_{2}\right]$ to be $-32.2 \mathrm{kcal} / \mathrm{mol}$ at 0 K and $-41.3 \mathrm{kcal} /$ mol at $298 \mathrm{~K} .{ }^{42}$ As in the $\mathbf{A B}$ monomer case, ${ }^{6}$ the error bar from the total atomization energies using $\operatorname{CCSD}(\mathrm{T}) / \mathrm{CBS}$ energies and other corrections is expected to be $\pm 1.0 \mathrm{kcal} / \mathrm{mol}$ for the heat of formation of the dimer. In conjunction with the
previously calculated heats of formation of the $\mathbf{A B}$ monomer ${ }^{6}$ $\Delta H_{\mathrm{f}}\left(\mathrm{BH}_{3} \mathrm{NH}_{3}\right)=-9.1 \pm 1.0 \mathrm{kcal} / \mathrm{mol}$ at 0 K and $\Delta H_{\mathrm{f}}\left(\mathrm{BH}_{3}-\right.$ $\left.\mathrm{NH}_{3}\right)=-13.5 \pm 1.0 \mathrm{kcal} / \mathrm{mol}$ at 298 K , the complexation energy of $\operatorname{dim}$ is $-14.0 \mathrm{kcal} / \mathrm{mol}$ at 0 K and $-14.3 \mathrm{kcal} / \mathrm{mol}$ at 298 K. Previous B3LYP ${ }^{31}$ and MP2 ${ }^{37}$ calculations predicted a complexation energy of about $-11 \mathrm{kcal} / \mathrm{mol}$ for $\mathbf{d i m}$. As shown in Table 2, the complexation energy of $\operatorname{dim}$ is calculated to be $-14.0,-13.7$, and $-13.5 \mathrm{kcal} / \mathrm{mol}$ using the $\operatorname{CCSD}(\mathrm{T})$ method with the aVTZ, aVQZ, and estimated CBS basis sets, respectively, including ZPE corrections. The result of $-13.5 \mathrm{kcal} /$ mol at the $\operatorname{CCSD}(\mathrm{T}) / \mathrm{CBS}+$ ZPE level differs by $0.5 \mathrm{kcal} / \mathrm{mol}$ from that of $-14.0 \mathrm{kcal} / \mathrm{mol}$ obtained previously ${ }^{42}$ This difference arises from the core-valence correlation and scalar relativistic corrections that were included in the previous work. Our current aVTZ value of $-14.0 \mathrm{kcal} / \mathrm{mol}$ for the complexation energy of the dimer dim is thus in good agreement with the previous higher-level calculations. ${ }^{42}$

Two isomers of dim with ion pair structures were also located. Selected MP2/aVTZ geometry parameters are also given in Figure 1. DADB contains a $\left[\mathrm{BH}_{4}^{-}\right]$anion whereas IonP-N bears a $\left[\mathrm{NH}_{4}{ }^{+}\right]$cation. Both isomers exhibit $C_{s}$ symmetry with the $\mathrm{B}-\mathrm{H}$ and $\mathrm{N}-\mathrm{H}$ bonds interacting strongly with each other. In DADB, the interion (B)H-H(N) distance of $1.539 \AA$ is relatively short, much shorter than the $\mathrm{H} \cdots \mathrm{H}$ distance of 1.986 $\AA$ in neutral dim. Similarly, the interion (N)H-H(B) distance in IonP-N is $1.535 \AA$. These geometric results show strong ionic interactions between the anion and the cation in the ion pairs.

We calculated the heats of formation of both ion pairs in the gas phase by using the total atomization energies computed at the $\operatorname{CCSD}(\mathrm{T}) / \mathrm{CBS}$ level plus additional corrections for the ZPE, core-valence effects, ${ }^{43}$ scalar relativistic effects, ${ }^{44}$ and atomic spin orbit effects ${ }^{45}$ following our prior work. ${ }^{6,7,13}$ The raw data for the atomization energies are given in Table S-2 of the Supporting Information. The atomization energy was converted to a heat of formation by using the 0 K atomic heats of formation of $\Delta H_{\mathrm{f}}{ }^{\circ}(\mathrm{N})=112.53 \pm 0.02 \mathrm{kcal} \mathrm{mol}^{-1}, \Delta H_{\mathrm{f}}{ }^{\circ}(\mathrm{B})=136.2 \pm$ $0.2 \mathrm{kcal} \mathrm{mol}^{-1}$, and $\Delta H_{\mathrm{f}}{ }^{\circ}(\mathrm{H})=51.63 \pm 0.001 \mathrm{kcal} \mathrm{mol}^{-1} .46,47$ Heats of formation at 298 K were obtained following the procedure of Curtiss et al. ${ }^{48} \mathrm{We}$ obtain $\Delta H_{\mathrm{f}}[\mathbf{D A D B}]=-21.6$ and $-31.0 \mathrm{kcal} / \mathrm{mol}$ at 0 and 298 K , respectively, and $\Delta H_{\mathrm{f}}[\mathbf{I}$ onP-N $]=-15.8$ and $-25.4 \mathrm{kcal} / \mathrm{mol}$ at 0 and 298 K . Accordingly, the DADB ion pair is found to be $10.6 \mathrm{kcal} / \mathrm{mol}$ less stable than dim. The IonP-N ion pair is even less stable, $16.4 \mathrm{kcal} / \mathrm{mol}$ above dim. We compile in Table 3 our calculated values at both 0 and 298 K , together with other relevant literature results. ${ }^{6,42}$

In the solid state, the ion pairs are expected to exist as ionic salts, and the lattice energies resulting from the salts could stabilize them differently. The corresponding volume for the salt $\left[\mathrm{BH}_{4}^{-}\right]\left[\mathrm{NH}_{3} \mathrm{BH}_{2} \mathrm{NH}_{3}{ }^{+}\right]$is $V_{\mathrm{m}}=0.1452 \mathrm{~nm}^{3}$ obtained from the sum of the cation, $V_{+}\left(\mathrm{NH}_{3} \mathrm{BH}_{2} \mathrm{NH}_{3}{ }^{+}\right)=0.0792 \mathrm{~nm}^{3}$, and the anion, $V_{-}\left(\mathrm{BH}_{4}^{-}\right)=0.066 \mathrm{~nm}^{3}$. The lattice energy of the crystal salt from Coulombic interaction between both $\left[\mathrm{BH}_{4}{ }^{-}\right]$ and $\left[\mathrm{NH}_{3} \mathrm{BH}_{2} \mathrm{NH}_{3}{ }^{+}\right]$ions is calculated to be $131.5 \mathrm{kcal} / \mathrm{mol}$ at 0 K , using eq 2, and $133.9 \mathrm{kcal} / \mathrm{mol}$ at 298 K , using eq 3 . The heats of formation of the ions in the gas phase are $\Delta H_{0 K}\left(\mathrm{BH}_{4}^{-}\right)$ $=-11.6 \mathrm{kcal} / \mathrm{mol}, \Delta H_{298 \mathrm{~K}}\left(\mathrm{BH}_{4}{ }^{-}\right)=-13.5 \mathrm{kcal} / \mathrm{mol},{ }^{6} \Delta H_{0 K}\left(\mathrm{NH}_{3}{ }^{-}\right.$ $\left.\mathrm{BH}_{2} \mathrm{NH}_{3}{ }^{+}\right)=105.9 \mathrm{kcal} / \mathrm{mol}$, and $\Delta H_{298 \mathrm{~K}}\left(\mathrm{NH}_{3} \mathrm{BH}_{2} \mathrm{NH}_{3}{ }^{+}\right)=$ $99.2 \mathrm{kcal} / \mathrm{mol}$. The heat of formation of $\mathrm{NH}_{3} \mathrm{BH}_{2} \mathrm{NH}_{3}{ }^{+}$was calculated from total atomization energies at the $\operatorname{CCSD}(\mathrm{T}) / \mathrm{CBS}$ level with the additional corrections described above. The zeropoint corrections were obtained from MP2/aVTZ harmonic vibrational frequencies and scaled by a factor of 0.9684 obtained as described above. The various components are given in Table

S-2 of the Supporting Information. From a Born-Haber cycle, the heat of formation of the salt, $\left[\mathrm{BH}_{4}^{-}\right]\left[\mathrm{NH}_{3} \mathrm{BH}_{2} \mathrm{NH}_{3}{ }^{+}\right]$, is predicted to be $-37.2 \mathrm{kcal} / \mathrm{mol}$ at 0 K and $-48.2 \mathrm{kcal} / \mathrm{mol}$ at 298 K with an estimated error of $\pm 5 \mathrm{kcal} / \mathrm{mol}$ due mostly to the simple model for the solid-state energy. The gas-phase binding energy of the ion pair from the heats of formation is $116.9 \mathrm{kcal} / \mathrm{mol}$ at 0 K , showing that all but $\sim 15 \mathrm{kcal} / \mathrm{mol}$ of the lattice energy is due to the interaction of a single ion pair.

For the second salt, $\left[\mathrm{NH}_{4}^{+}\right]\left[\mathrm{BH}_{3} \mathrm{NH}_{2} \mathrm{BH}_{3}{ }^{-}\right]$, the corresponding volume is $V_{\mathrm{m}}=0.1324 \mathrm{~nm}^{3}$ obtained from the sum of the cation, $V_{+}\left(\mathrm{NH}_{4}^{+}\right)=0.0210 \mathrm{~nm}^{3}$, and the anion, $V_{-}\left(\mathrm{BH}_{3} \mathrm{NH}_{2} \mathrm{BH}_{3}{ }^{-}\right)$ $=0.1114 \mathrm{~nm}^{3}$. The lattice energy is calculated to be $134.8 \mathrm{kcal} /$ mol at 0 K and $137.2 \mathrm{kcal} / \mathrm{mol}$ at 298 K . The heats of formation of the ions in the gas phase are $\Delta H_{0 \mathrm{KK}}\left(\mathrm{NH}_{4}{ }^{+}\right)=153.6 \mathrm{kcal} /$ $\mathrm{mol}, \Delta H_{298 \mathrm{~K}}\left(\mathrm{NH}_{4}{ }^{+}\right)=150.9 \mathrm{kcal} / \mathrm{mol},{ }^{6} \Delta H_{0 \mathrm{~K}}\left(\mathrm{BH}_{3} \mathrm{NH}_{2} \mathrm{BH}_{3}{ }^{-}\right)$ $=-49.4 \mathrm{kcal} / \mathrm{mol}$, and $\Delta H_{298 \mathrm{~K}}\left(\mathrm{BH}_{3} \mathrm{NH}_{2} \mathrm{BH}_{3}{ }^{-}\right)=-55.3 \mathrm{kcal} /$ mol. The heat of formation of $\mathrm{BH}_{3} \mathrm{NH}_{2} \mathrm{BH}_{3}{ }^{-}$was calculated from total atomization energies at the $\operatorname{CCSD}(\mathrm{T}) / \mathrm{CBS}$ level with the additional corrections described above. The various components are given in Table S-2 of the Supporting Information. The heat of formation of the salt, $\left[\mathrm{NH}_{4}{ }^{+}\right]\left[\mathrm{BH}_{3} \mathrm{NH}_{2} \mathrm{BH}_{3}{ }^{-}\right]$, is estimated to be $-30.7 \mathrm{kcal} / \mathrm{mol}$ at 0 K and $-41.6 \mathrm{kcal} / \mathrm{mol}$ at 298 K with an estimated error of $\pm 5 \mathrm{kcal} / \mathrm{mol}$. The difference in the heats of formation of the salts at 0 K is $6.5 \mathrm{kcal} / \mathrm{mol}$ favoring $\left[\mathrm{BH}_{4}^{-}\right]\left[\mathrm{NH}_{3} \mathrm{BH}_{2} \mathrm{NH}_{3}{ }^{+}\right]$, very similar to the difference of $5.8 \mathrm{kcal} / \mathrm{mol}$ for the gas-phase ion pairs. The ionic binding energy in the gas phase between the two ions is $120.0 \mathrm{kcal} / \mathrm{mol}$ at 0 K , and, again, all but $\sim 15 \mathrm{kcal} / \mathrm{mol}$ of the lattice energy is due to the interaction of a single ion pair.

From the gas-phase heats of formation of $\mathrm{BH}_{3} \mathrm{NH}_{2} \mathrm{BH}_{2} \mathrm{NH}_{3}$, linear trans and coiled (gauche) structures, ${ }^{42}$ we can calculate the following $\mathrm{H}_{2}$ elimination energies from the salts in kcal/ mol at 298 K

$$
\begin{array}{r}
{\left[\mathrm{BH}_{4}^{-}\right]\left[\mathrm{NH}_{3} \mathrm{BH}_{2} \mathrm{NH}_{3}^{+}\right] \rightarrow \mathrm{BH}_{3} \mathrm{NH}_{2} \mathrm{BH}_{2} \mathrm{NH}_{3} \text { (linear) }+} \\
\mathrm{H}_{2} \quad \Delta H_{\mathrm{rxn}}=12.9 \\
{\left[\mathrm{BH}_{4}^{-}\right]\left[\mathrm{NH}_{3} \mathrm{BH}_{2} \mathrm{NH}_{3}^{+}\right] \rightarrow \mathrm{BH}_{3} \mathrm{NH}_{2} \mathrm{BH}_{2} \mathrm{NH}_{3}(\text { coiled })+} \\
\mathrm{H}_{2} \quad \Delta H_{\mathrm{rxn}}=0.6 \tag{4b}
\end{array}
$$

$\left[\mathrm{NH}_{4}^{+}\right]\left[\mathrm{BH}_{3} \mathrm{NH}_{2} \mathrm{BH}_{3}^{-}\right] \rightarrow \mathrm{BH}_{3} \mathrm{NH}_{2} \mathrm{BH}_{2} \mathrm{NH}_{3}$ (linear) +
$\mathrm{H}_{2} \Delta H_{\mathrm{rxn}}=6.3$

$$
\begin{array}{r}
{\left[\mathrm{NH}_{4}^{+}\right]\left[\mathrm{BH}_{3} \mathrm{NH}_{2} \mathrm{BH}_{3}^{-}\right] \rightarrow \mathrm{BH}_{3} \mathrm{NH}_{2} \mathrm{BH}_{2} \mathrm{NH}_{3}(\text { coiled })+} \\
\mathrm{H}_{2} \Delta H_{\mathrm{rxn}}=-6.0 \tag{5b}
\end{array}
$$

The reaction 5b giving the stable coiled structure is exothermic, whereas formation of the linear structure is endothermic. Reaction 4b is very close to thermoneutral.

Reaction Pathways For $\mathbf{H}_{2}$ Release from Ammonia Borane Dimer. We first consider the energy barriers for the production of $\mathrm{H}_{2}$ molecules from dim. Geometries of the relevant TS's are displayed in Figures 2 and 4, and the potential energy diagrams constructed at the $\operatorname{CCSD}(\mathrm{T})$ level with different basis sets are illustrated in Figures 3 and 5. Results for the relative energy of some important structures as a function of basis sets are given in Table 2, whereas relative energies obtained with both aVDZ and aVTZ basis sets are listed in Table 4. Total energies, ZPEs, and Cartesian coordinates of all of the optimized structures are listed in Tables S-3 and S-4 of the Supporting Information. Geometry parameters of three additional TS's denoted as dim-ts6, dim-ts7, and dim-ts8 for $\mathrm{H}_{2}$ release from dim were also optimized and shown in Figure S-1 of the

TABLE 4: $\operatorname{CCSD}(T)$ Relative Energies Including ZPE Corrections ( $\mathrm{kcal} / \mathrm{mol}$ ) of Equilibrium and Transition State Structures for $\mathbf{H}_{2}$ Release from Ammonia Borane Dimer ${ }^{a}$

| structures | aVDZ | aVTZ |
| :--- | ---: | ---: |
| 2AB | 0.0 | 0.0 |
| dim | -14.2 | -14.0 |
| dim-ts1 | 36.3 | 36.0 |
| dim-ts2 | 31.2 | 30.5 |
| dim-ts3 | 32.8 | 32.2 |
| dim-ts4 | 34.3 | 33.9 |
| dim-ts5 | 34.9 | 34.1 |
| dim-p1 $+2 \mathrm{H}_{2}$ | -13.8 | -15.4 |
| 2( $\left.\mathrm{BH}_{2} \mathrm{NH}_{2}\right)+2 \mathrm{H}_{2}$ | -10.7 | -12.6 |
| dim-p2 $+\mathrm{H}_{2}$ | -11.5 | -12.2 |
| AB $+\mathrm{BH}_{2} \mathrm{NH}_{2}+\mathrm{H}_{2}$ | -5.4 | -6.3 |

${ }^{a}$ Calculated at the $\operatorname{CCSD}(\mathrm{T})$ level using MP2/aVTZ-optimized geometries and scaled MP2/aVDZ ZPE corrections.

Supporting Information. These TS's are however located at much higher energies and therefore not discussed further in the following section.
dim-ts 1 is the TS for generating two $\mathrm{H}_{2}$ molecules with each monomer contributing one $\mathrm{H}_{2}$ molecule (Figure 2). Formation of each $\mathrm{H}_{2}$ occurs within a $\mathrm{B}-\mathrm{H} \cdots \mathrm{H}-\mathrm{N}$ dim framework, involving a short $\mathrm{H}-\mathrm{H}$ distance of $0.863 \AA$, close to the bond distance of $0.74 \AA$ in $\mathrm{H}_{2}$. The reaction proceeds by an eightmembered ring TS. The $\mathrm{N}-\mathrm{H}$ distance of $1.594 \AA$ is proportionally more stretched than the $\mathrm{B}-\mathrm{H}$ bond $(1.396 \AA)$. The $\mathrm{B}-$ H and $\mathrm{N}-\mathrm{H}$ bonds distances of 1.197 and $1.011 \AA$ in the molecule are shorter than that of the $\mathrm{B}-\mathrm{H} \cdots \mathrm{H}-\mathrm{N}$ dihydrogen bond. A high barrier of $49.4 \mathrm{kcal} / \mathrm{mol}$ is found (Figure 3, CCSD(T)/CBS values). The two aminoborane molecules $\left(\mathrm{BH}_{2} \mathrm{NH}_{2}\right)$ also form a weak complex dim-p1 in the exit channel with a complexation energy of $-2.4 \mathrm{kcal} / \mathrm{mol}$. The product $2\left(\mathrm{BH}_{2} \mathrm{NH}_{2}\right)$ $+2 \mathrm{H}_{2}$ lies $0.5 \mathrm{kcal} / \mathrm{mol}$ below dim, which is close to a thermoneutral reaction.

As shown in Table 2, the energy barriers converge quickly with respect to the basis set. The values obtained at the CCSD(T)/aVTZ and $\operatorname{CCSD}(\mathrm{T}) / \mathrm{aVQZ}$ levels for dim differ only by, at most, 0.5 and $0.2 \mathrm{kcal} / \mathrm{mol}$, respectively, from the CCSD(T)/CBS counterparts. This is consistent with what we have found previously. ${ }^{13,26}$ These results show that use of the computationally less-demanding $\operatorname{CCSD}(\mathrm{T}) /$ aVTZ level can be used to construct the potential energy surfaces for these systems, with an expected accuracy of $\pm 1.0 \mathrm{kcal} / \mathrm{mol}$ for the energy barriers as compared to the CBS limit.

In addition to the dim-ts1 described above, we located four other TS's for release of one $\mathrm{H}_{2}$ molecule from dim. Figure 4 shows critical MP2/aVTZ geometry parameters of these TS's, and Figure 5 summarizes the reaction pathway energetics.
dim-ts2 is basically an $\mathrm{H}_{2}$ elimination from one $\mathbf{A B}$ monomer ${ }^{13}$ with a TS that is complexed with the second monomer by $\mathrm{H}-\mathrm{H}$ interactions from the side opposite to the departing $\mathrm{H}_{2}$. The intermolecular $\mathrm{H}-\mathrm{H}$ distances range from 2.069 to $2.305 \AA$, and a $\mathrm{B}-\mathrm{H} \cdots \mathrm{H}-\mathrm{N}$ skeleton is present with bond distances of $1.211,2.069$, and $1.011 \AA$, respectively. The energy for this TS is $30.5 \mathrm{kcal} / \mathrm{mol}$ above the separated monomers and $44.5 \mathrm{kcal} / \mathrm{mol}$ above the dimer $\operatorname{dim}$ (values at the $\operatorname{CCSD}(\mathrm{T})$ / $\mathrm{aVTZ}+$ ZPE level, Figure 5). The product dim-p2 is characterized as a hydrogen-bonded complex between $\mathrm{BH}_{3} \mathrm{NH}_{3}$ and $\mathrm{BH}_{2} \mathrm{NH}_{2}$, with a complexation energy of $-5.9 \mathrm{kcal} / \mathrm{mol}$. The product dim-p2 $+\mathrm{H}_{2}$ is calculated to be $1.8 \mathrm{kcal} / \mathrm{mol}$ above dim.
dim-ts3 corresponds to a comparable process to that involving dim-ts2, but the intermolecular interaction mainly involves a type of $\mathrm{N} \cdots \mathrm{H}-\mathrm{N}$ hydrogen bond. The $\mathrm{N}-\mathrm{H}$ bond of the


Figure 6. An IRC pathway showing the formation of the product dim-p2 from dim-ts4, calculated using the B3LYP/6-311++G(d,p) level. IRC calculations starting from the TS were done with the default step size of $0.1 \mathrm{amu}^{1 / 2}$ bohr.
nonreacting monomer plays the role of the hydrogen-bond donor, with an intermolecular $\mathrm{N} \cdots \mathrm{H}$ distance of $2.240 \AA$. Parameters within the four-membered cycles involving the forming $\mathrm{H}_{2}$ in both dim-ts2 and dim-ts3 differ somewhat from each other, with $\mathrm{B}-\mathrm{N}$ distances varying up to $0.04 \AA$. The role of the second monomer is described in more detail below. The energy of dim-ts $\mathbf{3}$ is higher, by $1.7 \mathrm{kcal} / \mathrm{mol}$, than dim-ts $\mathbf{2}$. The process through dim-ts3 also gives rise to the same product dim-p2.

The reaction path proceeding through dim-ts4 shows another way of eliminating $\mathrm{H}_{2}$. The departing $\mathrm{H}_{2}$ is composed of one $\mathrm{H}(\mathrm{B})$ atom of the first monomer plus one $\mathrm{H}(\mathrm{N})$ atom of the second monomer. The strongest intermolecular interaction arises from a $\mathrm{N}-\mathrm{H} \cdots \mathrm{N}$ hydrogen bond involving quite short $\mathrm{N}-\mathrm{H}$ distances of 1.138 and $1.475 \AA$, so the H is partially transferred. In dim-ts4, both $\mathrm{BH}_{3}$ groups are trans with respect to the N $\mathrm{H} \cdots \mathrm{N}$ bond. The $\mathrm{B}-\mathrm{N}-\mathrm{H} \cdots \mathrm{N}-\mathrm{B}$ frame lies almost in the same plane. The $\mathrm{H}-\mathrm{H}$ distance of $0.811 \AA$ shows that the product $\mathrm{H}_{2}$ molecule is almost completely formed. dim-ts4 is close to the TS for $\mathrm{H}_{2}$ release from $\mathrm{BH}_{3} \mathrm{NH}_{3}$ assisted by a $\mathrm{NH}_{3}$ molecule that we previously reported. ${ }^{13}$ In this case, the second monomer plays the role of the bifunctional catalyst. The channel involving dim-ts4 has energy barriers of 33.9 and $47.9 \mathrm{kcal} / \mathrm{mol}$ with respect to two $\mathbf{A B}$ monomers and dim, respectively, to produce one $\mathrm{H}_{2}$. dim-ts4 is $1.7 \mathrm{kcal} / \mathrm{mol}$ higher in energy than dim-ts3 and could lead to the chain product DADB-p1 $\left(\mathrm{BH}_{3} \mathrm{NH}_{2} \mathrm{BH}_{2}-\right.$ $\mathrm{NH}_{3}$ ) involved in reactions 4 and 5 . Upon loss of $\mathrm{H}_{2}$, both tricoordinated B and N atoms could join to make a new $\mathrm{B}-\mathrm{N}$ bond. However, extensive IRC calculations demonstrate that an H transfer between two N atoms takes place yielding the complex dim-p2. The corresponding IRC path for product formation from dim-ts4 at the B3LYP/6-311++G(d,p) level is displayed as Figure 6.
dim-ts5 involves a similar kind of process to that proceeding through dim-ts4 with the main difference between dim-ts4 and dim-ts5 being that in the latter the $\mathrm{BH}_{3}(\mathrm{~N})$ group of the second monomer is nearly perpendicular to the $\mathrm{B}-\mathrm{N}-\mathrm{N}$ plane. The intermolecular distances in dim-ts5 are markedly longer than those in dim-ts4. dim-ts5 corresponds to the energetically highest-lying TS among the four TS's examined. It still is only $3.6 \mathrm{kcal} / \mathrm{mol}$ above dim-ts2. Again, dim-ts5 leads to the product dim-p2.

The energy profiles summarized in Figure 5 show that the energies of the five TS's do not differ substantially from each other. The difference between the lowest-lying dim-ts2 and the highest-lying dim-ts1 is $5.5 \mathrm{kcal} / \mathrm{mol}$. The catalytic effect of simple dihydrogen-bonded interactions via dim-ts2 and dimts3 is slightly more effective than the bifunctional hydrogen exchange via dim-ts4 and dim-ts5. Nevertheless, the smallest barrier is $44.5 \mathrm{kcal} / \mathrm{mol}$ above the dimer dim. This is much larger than the dimerization energy of $\mathbf{A B}$ of only $14.0 \mathrm{kcal} /$ mol (values at $\operatorname{CCSD}(\mathrm{T}) / \mathrm{aVTZ}+\mathrm{ZPE}$ ) and is above the dimerization energy plus breaking one $\mathbf{A B} \mathrm{B}-\mathrm{N}$ bond. Thus, the dimer dim is expected to undergo dissociation rather than $\mathrm{H}_{2}$ production under mild thermal conditions.

Formation of the Diammoniate of Diborane Ion Pair (DADB) and $\mathbf{H}_{2}$ Release from Two Ammonia Borane Molecules. We now consider the formation of DADB from three distinct starting points: $2 \mathrm{BH}_{3} \mathrm{NH}_{3}, \mathrm{~B}_{2} \mathrm{H}_{6}+2 \mathrm{NH}_{3}$, and $\mathrm{BH}_{3}-$ $\mathrm{NH}_{3}+\mathrm{NH}_{3}+\mathrm{BH}_{3}$. Selected geometry parameters of all related structures are given in Figure 7. These include four TS's for the formation of DADB (Figure 7a), three complexes formed from $\mathrm{BH}_{3} \mathrm{HBH}_{2} \mathrm{NH}_{3}+\mathrm{NH}_{3}$, with different orientations of the $\mathrm{NH}_{3}$ (Figure 7b), and also three TS's for $\mathrm{H}_{2}$ release from intermediates (Figure 7c). Figure 8 schematically illustrates the reaction pathways leading to DADB formation and to $\mathrm{H}_{2}$ formation.

As shown in Figure 8, all three reactant systems 2AB, $\mathbf{B}_{2} \mathrm{H}_{6}$ $+2 \mathrm{NH}_{3}$ (denoted as $\mathbf{D B}+2 \mathrm{~A}$ ), and $\mathrm{BH}_{3} \mathrm{NH}_{3}+\mathrm{NH}_{3}+\mathrm{BH}_{3}$ (denoted as $\mathbf{A B}+\mathrm{A}+\mathrm{B}$ ) undergo additions and rearrangements converging to the same adduct BBA-com +A in the first step. BBA-com is a key common intermediate in these processes and is the adduct of $\mathbf{A B}$ with $\mathrm{BH}_{3}$, which was studied in detail in our previous paper. ${ }^{13}$ Thus, condensation of $\mathbf{A B}+\mathrm{A}+\mathrm{B}$ invariably leads to BBA-com $+B$ without an activation energy barrier.

BBA-com + B is calculated to be $8.1 \mathrm{kcal} / \mathrm{mol}$ above the two separated monomers $\mathbf{2 A B}$ and substantially below $\mathbf{A B}+$ $\mathrm{B}+\mathrm{A}$. Addition from $\mathbf{2 A B}$ involves an energy barrier of 12.4 $\mathrm{kcal} / \mathrm{mol}$ via TS 2AB-ts1. The TS involves the attack of the B atom in one $\mathbf{A B}$ on an H atom of the $\mathrm{BH}_{3}$ in another $\mathbf{A B}$. Formation of a bridge $\mathrm{B}-\mathrm{H}-\mathrm{B}$ bond induces breakage of the $\mathrm{N}-\mathrm{B}$ bond in the attacking $\mathbf{A B}$ to eliminate $\mathrm{NH}_{3}$. The BH distances of 1.700 and $1.229 \AA$ in 2 AB-ts 1 show a forming


Figure 7. Selected MP2/aVTZ geometry parameters of four different TS's for formation of $\mathbf{D A D B}$ from $\mathbf{2 A B}, \mathbf{D B}+2 \mathrm{~A}$, and $\mathbf{A B}+\mathbf{A}+\mathrm{B}$. Selected MP2/aVTZ geometry parameters of three different complexes for formation of $\mathbf{D A D B}$ from $\mathbf{2 A B}, \mathbf{D B}+2 \mathrm{~A}$, and $\mathbf{A B}+\mathrm{A}+\mathrm{B} . \operatorname{Selected}$ MP2/aVTZ geometry parameters of three different TS's for $H_{2}$ release from $\mathbf{2 A B}, \mathbf{D B}+2 \mathrm{~A}$, and $\mathbf{A B}+\mathrm{A}+\mathrm{B}$. Bond lengths are given in $\AA$, and bond angles in degrees.
$B-H-B$ bond. The stretching of the $N-B$ bond in the second AB by $\sim 0.8 \AA$ to $2.415 \AA$ is consistent with breaking a B-N bond (cf. Figure 7a).

The barrier-free condensation of diborane (DB) with one $\mathrm{NH}_{3}$ molecule forms the complex DBA-com, which has a complexation energy of $-5.1 \mathrm{kcal} / \mathrm{mol}$ and is $12.6 \mathrm{kcal} / \mathrm{mol}$ above $\mathbf{2 A B}$. This weak complex is characterized by long distances between N and a bridge BH bond (not shown in Figure 7b). Conversion of DBA-com to BBA-com occurs with an energy barrier of $7.8 \mathrm{kcal} / \mathrm{mol}$ through TS DBA-ts, which lies $20.4 \mathrm{kcal} / \mathrm{mol}$ above $\mathbf{2 A B}$. DBA-ts is formed by a nearly perpendicular attack of N of a free $\mathrm{NH}_{3}$ to a B center of $\mathbf{D B}\left(\angle \mathrm{NBH}\right.$ angle $=78.4^{\circ}$ and $\mathrm{N}-\mathrm{B}=2.042 \AA$ ). When $\mathrm{NH}_{3}$ approaches the diborane plane defined by $\mathrm{BH}_{\mathrm{b}} \mathrm{BH}_{\mathrm{b}}$, the attack on one $\mathrm{B}-\mathrm{H}-\mathrm{B}$ bond induces breaking of the other $\mathrm{B}-\mathrm{H}-\mathrm{B}$ bond. The broken $\mathrm{B}-\mathrm{H}$ distance is $1.604 \AA$ in DBA-ts (Figure 7a). The identity of both TS's was confirmed by IRC calculations at the B3LYP level with the $6-311++G(d, p)$ basis set. ${ }^{49}$

At this point, we have described the formation of BBA-com + A as separated species that represent another asymptotic region. This asymptote provides a common entry point to an interesting region of the potential energy surface. When BBAcom and A are allowed to interact, they exothermically form without a barrier a variety of hydrogen-bonded complexes depending on the orientation of the $\mathrm{NH}_{3}$. Figure 7 shows three
complexes 2AB-com1, 2AB-com2, and 2AB-com3 with different formation energies and related to DADB or $\mathrm{H}_{2}$ formation. Of these complexes, $\mathbf{2 A B}-\mathbf{c o m} 1$ is located at the lowest-energy position, $1.0 \mathrm{kcal} / \mathrm{mol}$ below $\mathbf{2 A B}$, and has an intermolecular $\mathrm{N}-\mathrm{H}$ bond between the N of the free $\mathrm{NH}_{3}$ and an H of the $\mathrm{NH}_{3}$ in BBA-com. 2AB-com1 is characterized by the largest complexation energy of $-9.1 \mathrm{kcal} / \mathrm{mol}$ with respect to BBA$\mathbf{c o m}+\mathrm{A}$. The small energy difference between $\mathbf{2 A B}$ and $\mathbf{2 A B}$ com1 means that the energy lost by breaking the $\mathrm{B}-\mathrm{N}$ bond in an $\mathbf{A B}$ monomer is fully recovered by the new $\mathrm{B}-\mathrm{H}-\mathrm{B}$ and $\mathrm{N}-\mathrm{H}-\mathrm{N}$ bonds in the complex.

2AB-com 2 is a slightly weaker complex, formed by the same type of $\mathrm{N}-\mathrm{H}-\mathrm{N}$ interaction as in 2AB-com1, but with a different of $\mathrm{NH}_{3}$ orientation. In this case, the $\mathrm{NH}_{3}$ is connected to one H atom of the bonded $\mathrm{NH}_{3}$ group and is in a trans position with respect to the $\mathrm{BH}_{3}$ of BBA-com complexed to the H from the $\mathrm{BH}_{3} \mathrm{NH}_{3}$ moiety. $\mathbf{2 A B}$-com2 2 is only $0.7 \mathrm{kcal} / \mathrm{mol}$ less stable than $\mathbf{2 A B}$-com1 and is $0.3 \mathrm{kcal} / \mathrm{mol}$ below $\mathbf{2 A B}$. There is a spectrum of complexes between BBA-com and $\mathrm{NH}_{3}$ that differ from each other by rotation of the $\mathrm{NH}_{3}$. These are however less stable than the two complexes just discussed.

2AB-com3 has four $\mathrm{N}-\mathrm{H}-\mathrm{H}-\mathrm{B}$ interactions between the $\mathrm{N}-\mathrm{H}$ bonds of the free $\mathrm{NH}_{3}$ and $\mathrm{B}-\mathrm{H}$ bonds belonging to BBAcom. The $\mathrm{H}-\mathrm{H}$ bond distances are long, $\sim 3.0 \AA$, and the complex is weak with a complexation energy of only $-1.1 \mathrm{kcal} /$


Figure 8. Schematic energy profiles illustrating the different reaction pathways for formation of $\mathbf{D A D B}$ from $\mathbf{2 A B}, \mathbf{D B}+2 \mathrm{~A}$, and $\mathbf{A B}+\mathrm{A}+\mathrm{B}$. Relative energies given in $\mathrm{kcal} / \mathrm{mol}$ were obtained from $\operatorname{CCSD}(\mathrm{T}) / \mathrm{aVTZ}+$ ZPE calculations.
mol with respect to separated BBA-com + A fragments. 2ABcom3 is $7.0 \mathrm{kcal} / \mathrm{mol}$ above $\mathbf{2 A B}$ and $8.0 \mathrm{kcal} / \mathrm{mol}$ above $\mathbf{2 A B}-$ com1.

As shown in Figure 8, DADB can be formed from unimolecular rearrangements of $2 \mathrm{AB}-\operatorname{com} 2$ and $2 \mathrm{AB}-\operatorname{com} 3$ but not from 2 AB -com1. 2 AB -com 2 is connected to DADB by an energy barrier of $28.1 \mathrm{kcal} / \mathrm{mol}$ via TS 2AB-ts3. This is the sole TS from 2AB-com2 and is located at $27.8 \mathrm{kcal} / \mathrm{mol}$ above $\mathbf{2 A B}$. Conversion of $\mathbf{2 A B}$-com3 to DADB has a larger energy barrier of $31.2 \mathrm{kcal} / \mathrm{mol}$ through TS 2AB-ts5 $(38.2 \mathrm{kcal} / \mathrm{mol}$ above $\mathbf{2 A B}$ ). Geometries of these TS's are shown in Figure 7.

DADB can also be formed from $\mathbf{2 A B}$-com 3 by overcoming an energy barrier of $31.2 \mathrm{kcal} / \mathrm{mol}$ at TS 2AB-ts5 ( $38.2 \mathrm{kcal} /$ mol above $\mathbf{2 A B}$ ). The higher energy of the latter TS, relative to $\mathbf{2 A B}-\mathbf{t s 3}$, is due to different factors. A simultaneous breaking of not only the bridged $B-H-B$ but also four $N-H-H-B$ bonds, even though they are weak, occurs. In 2AB-ts5, the new BN and BH bonds and, thus, both cationic and anionic moieties have essentially been formed. Such a "late" transition state is usually associated with a higher barrier. However, formation of DADB is not the favored rearrangement of 2 AB -com3. Indeed it undergoes preferentially loss of $\mathrm{H}_{2}$ with a lower-energy barrier of $23.7 \mathrm{kcal} / \mathrm{mol}$ via TS 2AB-ts4 (Figure 7, $30.7 \mathrm{kcal} /$
mol above $\mathbf{2 A B}$ ), producing cyclic $\mathbf{2 A B}-\mathbf{p 1}$. The latter TS $\mathbf{2 A B}-$ ts4 is basically composed of the TS for $\mathrm{H}_{2}$ loss from BBAcom, reported in our previous paper, ${ }^{13}$ with complexation by an additional $\mathrm{NH}_{3}$. A difference of $7.5 \mathrm{kcal} / \mathrm{mol}$ between the energies of both TS's clearly points out the dominance of the $\mathrm{H}_{2}$ elimination over DADB formation from 2AB-com3.

We also found another pathway for $\mathrm{H}_{2}$ release from 2ABcom1 via TS 2AB-ts2 in which one $\mathrm{H}_{2}$ can be generated by the same mechanism as in BBA-com ${ }^{13}$ and again complexed by $\mathrm{NH}_{3}$ but on the amine side. This TS is close to $2 \mathrm{AB}-\mathrm{ts} 4$, but they differ from each other by the position of the $\mathrm{NH}_{3}$ and different hydrogen-bond interactions. 2AB-ts2 is lower in energy than 2AB-ts $\mathbf{4}$ by $4.9 \mathrm{kcal} / \mathrm{mol}$; it is 25.8 and $26.8 \mathrm{kcal} / \mathrm{mol}$ above $\mathbf{2 A B}$ and $2 \mathrm{AB}-c o m 1$, respectively. In this case, the products $\mathbf{2 A B}-\mathbf{p} 2+\mathrm{NH} 3+\mathrm{H}_{2}$ are only $0.2 \mathrm{kcal} / \mathrm{mol}$ above $\mathbf{2 A B}$.

From the $\mathbf{A B}+\mathrm{A}+\mathrm{B}$ asymptote, the interaction between $\mathbf{A B}$ and $\mathrm{NH}_{3}$ leads directly to the complex BAA-com, which was also studied in our previous study. ${ }^{13}$ BAA-com is characterized by an $\mathrm{N}-\mathrm{H}$ hydrogen-bond distance of $1.99 \AA$, and additional stability is provided by a weaker $\mathrm{B}-\mathrm{H}-\mathrm{N}$ interaction. (The $\mathrm{B}-\mathrm{H}\left(\mathrm{NH}_{3}\right)$ bond is $2.68 \AA$.) In this complex, $\mathrm{NH}_{3}$ plays the role of a hydrogen-bond acceptor, and $\mathbf{A B}$ is a hydrogenbond donor. The energy of BAA-com $+\mathrm{BH}_{3}$ is $17.9 \mathrm{kcal} / \mathrm{mol}$


Figure 9. Selected MP2/aVTZ geometry parameters of five different TS's for $H_{2}$ release from DADB and IonP-N. Bond lengths are given in $\AA$, and bond angles in degrees.
above $\mathbf{2 A B}$ and $7.4 \mathrm{kcal} /$ mol below $\mathbf{A B}+\mathrm{A}+\mathrm{B}$. With the assistance of $\mathrm{BH}_{3}$, one $\mathrm{H}_{2}$ molecule can be generated from BAAcom + B through TS 2AB-ts6 (Figure 7c) with an energy barrier of $20.6 \mathrm{kcal} / \mathrm{mol}$ ( $\mathbf{2 A B}-\mathrm{ts} 6$ is $38.5 \mathrm{kcal} / \mathrm{mol}$ above $\mathbf{2 A B}$ ). This TS results from an attack of the B of free $\mathrm{BH}_{3}$ to the N of $\mathrm{NH}_{3}$ in BAA-com, and the two departing H atoms are bonded at a distance of $1.134 \AA$. One H atom is released from $\mathrm{NH}_{3}$ in BBAcom and the other from free $\mathrm{BH}_{3}$. IRC calculations showed that the products of this TS are also $\mathbf{2 A B} \mathbf{- p} \mathbf{1}+\mathbf{A}+\mathrm{H}_{2}$ (as from 2AB-com3). 2AB-ts6 corresponds to a TS for $\mathrm{H}_{2}$ loss of a $\mathbf{A B}$ monomer interacting simultaneously with $\mathrm{BH}_{3}$ and $\mathrm{NH}_{3}$. As shown in Figure 8, the reaction pathway involving BAA-com $+B$ is less favored with respect to the paths discussed above. 2AB-ts6 represents the highest-lying saddle point in this region of the potential energy surface.

In summary, the results in Figure 8 show that a reaction starting from two monomers $\mathbf{2 A B}$ yields the complexes $\mathbf{2 A B}$ com1 and 2 AB -com 2 as reactive intermediates. The complexes undergo subsequent $\mathrm{H}_{2}$ release and DADB formation reactions characterized by energy barriers of $25.8 \mathrm{kcal} / \mathrm{mol}$ via $\mathbf{2 A B}-\mathrm{ts} 2$ and $27.8 \mathrm{kcal} / \mathrm{mol}$ via $\mathbf{2 A B}-t s 3$, respectively. An energy
difference of $2 \mathrm{kcal} / \mathrm{mol}$ between both TS's is small and close to the expected accuracy of our calculations at the $\operatorname{CCSD}(\mathrm{T})$ / $\mathrm{aVTZ}+$ ZPE level. Other factors besides the simple barrier heights could play a role in the kinetics. Thus both channels, $\mathrm{H}_{2}$ release and DADB formation, should be competitive. Although energetically less favored, the BAA-com +B and 2AB-com3 systems also connect other regions of the potential energy surface, such as the channel giving rise to diborane $\left(\mathrm{B}_{2} \mathrm{H}_{6}\right)$ as a product. We have recently probed the reaction of diborane with ammonia, which follows similar reaction pathways. ${ }^{26,49}$

We now address the role of the second $\mathbf{A B}$ monomer in the various transformations. We showed ${ }^{13}$ that $\mathrm{B}-\mathrm{N}$ bond cleavage is favored by $\sim 10 \mathrm{kcal} / \mathrm{mol}$ over an $\mathrm{H}_{2}$ loss from one $\mathbf{A B}$ monomer. The role of the second monomer is to provide the $\mathrm{BH}_{3}$ and $\mathrm{NH}_{3}$ that can catalyze loss of $\mathrm{H}_{2}$ from the remaining monomer or to assist in forming various complexes associated with breaking a $\mathrm{B}-\mathrm{N}$ bond. The reaction pathways related to the $\mathbf{2 A B}$ and $\mathbf{A B}+\mathrm{A}+\mathrm{B}$ systems given in Figure 8 show that the $\mathbf{A B}$ monomer can assist in breaking the $\mathrm{B}-\mathrm{N}$ bond as

TABLE 5: $\operatorname{CCSD}(T)$ Relative Energies ( $\mathbf{k c a l} / \mathrm{mol}$ ) of Transition State and Equilibrium Structures for Formation DADB and $H_{2}$ Release from the $2 \mathrm{AB}, \mathrm{DB}+2 \mathrm{~A}$, and $\mathrm{AB}+$ $\mathbf{A}+\mathbf{B}^{a}$

| structures | aVDZ | aVTZ |
| :--- | ---: | ---: |
| 2AB | 0.0 | 0.0 |
| DB $+2 A$ | 17.9 | 17.7 |
| AB + A + B | 24.0 | 25.3 |
| DBA-com + A | 12.2 | 12.6 |
| BAA-com + B | 16.5 | 17.9 |
| BBA-com + A | 8.0 | 8.1 |
| 2AB-com1 | -1.8 | -1.0 |
| 2AB-com2 | -0.9 | -0.3 |
| 2AB-com3 | 6.5 | 7.0 |
| DADB | -3.5 | -3.3 |
| 2AB-p1 + H2 | -14.8 | -16.7 |
| 2AB-p2 + A + $\mathrm{H}_{2}$ | 2.0 | 0.2 |
| DBA-ts + A | 19.9 | 20.4 |
| 2AB-ts1 | 11.0 | 12.4 |
| 2AB-ts2 | 25.0 | 25.8 |
| 2AB-ts3 | 26.7 | 27.8 |
| 2AB-ts4 | 30.3 | 30.7 |
| 2AB-ts5 | 37.2 | 38.2 |
| 2AB-ts6 | 38.5 | 38.5 |

${ }^{a}$ Calculated at the $\operatorname{CCSD}(\mathrm{T})$ level using MP2/aVTZ-optimized geometries and scaled MP2/aVDZ ZPE corrections.

TABLE 6: $\operatorname{CCSD}(\mathrm{T})$ Relative Energies ( $\mathrm{kcal} / \mathrm{mol}$ ) of Transition State and Equilibrium Structures for $\mathbf{H}_{2}$ Release from the Ion Pair, $\left[\mathrm{BH}_{4}{ }^{+}\right]\left[\mathrm{NH}_{3} \mathbf{B H}_{2} \mathbf{N H}_{3}{ }^{-}\right]$, DADB ${ }^{a}$

| structures | aVDZ | aVTZ |
| :--- | ---: | ---: |
| 2AB | 0.0 | 0.0 |
| DADB | -3.5 | -3.3 |
| DADB-p1 $+\mathrm{H}_{2}$ | -7.8 | -9.3 |
| DADB-p2 $+\mathrm{H}_{2}$ | 8.1 | 7.4 |
| DADB-p3 $+\mathrm{H}_{2}$ | -19.1 | -20.8 |
| DADB-ts1 | 17.2 | 17.2 |
| DADB-ts2 | 48.3 | 47.8 |
| DADB-p1-ts $+\mathrm{H}_{2}$ | -5.9 | -7.3 |

${ }^{a}$ Calculated at the $\operatorname{CCSD}(\mathrm{T})$ level using MP2/aVTZ-optimized geometries and scaled MP2/aVDZ ZPE corrections.
in 2AB-ts1 or can break apart, forming $\mathrm{BH}_{3}$ and $\mathrm{NH}_{3}$ to catalyze the process as in 2AB-ts2 and 2AB-ts3.

In addition, it is useful to compare $\mathbf{2 A B}$-ts1 and $\mathbf{2 A B}$-ts3 (Figure 7), which only differ by the orientation of the $\mathrm{NH}_{3}$ relative to the effectively separated $\mathrm{BH}_{3}$. In $\mathbf{2 A B}$-ts1, the $\mathrm{NH}_{3}$ interacts with the noncovalently bonded $\mathrm{BH}_{3}$ and connects BBAcom to $2 \mathbf{A B}$. In 2AB-ts3, the $\mathrm{NH}_{3}$ approaches the $\mathrm{BH}_{3}$ of the monomer from the opposite side, leading to transfer of the H on the $\mathrm{BH}_{3}$ to form $\mathrm{BH}_{4}^{-}$and $\mathrm{NH}_{3} \mathrm{BH}_{2} \mathrm{NH}_{3}{ }^{+}$, the components of DADB. Thus, the orientation of the incoming $\mathrm{NH}_{3}$ controls which part of the potential energy surface is accessed.

Reaction Pathways for $\mathbf{H}_{\mathbf{2}}$ Release from Ion Pair Isomers. We now examine the production of $\mathrm{H}_{2}$ from the ion pairs DADB and IonP-N. Figure 9 displays geometrical parameters of five relevant TS's. Calculated relative energies obtained at the $\operatorname{CCSD}(\mathrm{T})$ level with the aVDZ and aVTZ basis sets are given in Tables 5, 6, and 7. The energy profiles at the $\operatorname{CCSD}(\mathrm{T})$ / aVTZ level are shown in Figures 10, 11, and 12. Total and ZPEs are tabulated in Table S-3 of the Supporting Information.

DADB-ts1 is the TS for the generation of one $\mathrm{H}_{2}$ molecule from DADB (Figure 9). For this reaction channel, results up to the $\operatorname{CCSD}(\mathrm{T}) / \mathrm{CBS}$ level are compared in Table 2 and illustrated in Figure 10. Again, the energies obtained using the aVTZ basis set are very similar to the CBS values. As hydrogen transfer occurs, the $\mathrm{BH}_{4}$ moiety rotates around a $\mathrm{NH}_{3}$ group to yield a


Figure 10. Schematic potential energy profiles in $\mathrm{kcal} / \mathrm{mol}$ showing the reaction pathway for $\mathrm{H}_{2}$ release from DADB, calculated with the $\operatorname{CCSD}(\mathrm{T})$ method using four basis sets including zero-point corrections. The energies are from top to bottom: aVDZ, aVTZ, aVQZ, and CBS.

TABLE 7: CCSD(T) Relative Energies ( $\mathrm{kcal} / \mathrm{mol}$ ) of Transition State and Equilibrium Structures for $\mathrm{H}_{2}$ Release from the Ion Pair, $\left[\mathrm{NH}_{4}{ }^{+}\right]\left[\mathrm{BH}_{3} \mathbf{N H}_{2} \mathrm{BH}_{3}{ }^{-}\right]$, IonP- ${ }^{a}$

| structures | aVDZ | aVTZ |
| :--- | :---: | :---: |
| 2AB | 0.0 | 0.0 |
| IonP-N | 2.3 | 2.7 |
| IonP-N $+\mathrm{H}_{2}$ | 14.6 | 14.2 |
| IonP-N-ts1 | 20.3 | 20.4 |
| IonP-N-ts2 | 59.0 | 58.7 |

${ }^{a}$ Calculated at the $\operatorname{CCSD}(\mathrm{T})$ level using MP2/aVTZ-optimized geometries and scaled MP2/aVDZ ZPE corrections.
trans configuration for DADB-ts1. IRC calculations confirmed the connection between DADB and the product DADB-p1. The $\mathrm{B}-\mathrm{H}$ bond of the $\mathrm{BH}_{4}^{-}$moiety only lengthens by $0.09 \AA$, whereas the $\mathrm{N}-\mathrm{H}$ bond is substantially longer by $0.37 \AA$. The forming $\mathrm{H}_{2}$ has a distance of $0.968 \AA$ at DADB-ts1, longer than the values found above. $\mathrm{H}_{2}$ production through this TS is of much lower energy than the one described above, with an energy barrier of only $20.1 \mathrm{kcal} / \mathrm{mol}$. The product generated from DADB is the chain molecule $\mathrm{NH}_{3} \mathrm{BH}_{2} \mathrm{NH}_{2} \mathrm{BH}_{3}$ DADB-p1. The reaction is slightly exothermic, $-6.5 \mathrm{kcal} / \mathrm{mol}$ with respect to DADB (CBS results, Figure 10). The trans chain (linear) DADB-p1 found by IRC calculations as the primary product does not correspond to the lowest-lying conformation of the molecule. As given in reactions 4 and 5, the gauche (coiled) conformer is $12.2 \mathrm{kcal} / \mathrm{mol}$ more stable than the trans conformer, due to the presence of stabilizing $\mathrm{B}-\mathrm{H} \cdots \mathrm{H}-\mathrm{N}$ interactions ${ }^{42}$ ( $11.5 \mathrm{kcal} / \mathrm{mol}$ at the $\operatorname{CCSD}(\mathrm{T}) / \mathrm{aVTZ}+\mathrm{ZPE}$ level, cf. Figure 11). The conversion of the trans DADB-p1 to the gauche form DADB-p3 described in Figure 11 is however a low-energy process, with a barrier of rotation of $2.0 \mathrm{kcal} / \mathrm{mol}$ around the central $\mathrm{B}-\mathrm{N}$ bond involving the TS DADB-p1-ts (Figure 9).

There is another pathway for generating $\mathrm{H}_{2}$ from DADB involving the TS DADB-ts2 leading to the ion pair DADB-p2. As shown in Figure 11, DADB-p2 is $16.7 \mathrm{kcal} / \mathrm{mol}$ less stable than DADB-p1, and DADB-ts2 lies $30.6 \mathrm{kcal} / \mathrm{mol}$ higher in energy thanDADB-ts1. Therefore, the ionpair $\left[\mathrm{BH}_{4}{ }^{-}\right]\left[\mathrm{NH}_{3} \mathrm{BH}_{2} \mathrm{NH}_{3}{ }^{+}\right]$ undergoes preferentially $\mathrm{H}_{2}$ eliminations overcoming an energy barrier of $20.5 \mathrm{kcal} / \mathrm{mol}$ via DADB-ts1 (Figure 11). This result is consistent with the experimental finding that DADB, gener-


Figure 11. Schematic energy profiles illustrating the different reaction pathways for $\mathrm{H}_{2}$ release from DADB. Relative energies in $\mathrm{kcal} / \mathrm{mol}$ were obtained from $\operatorname{CCSD}(\mathrm{T}) / \mathrm{aVTZ}+\mathrm{ZPE}$ calculations.
ated following a seeding of $(\mathbf{A B})_{2}$ dimer, releases $\mathrm{H}_{2}$ under mild temperatures. ${ }^{14}$

IonP-N-ts1 and IonP-N-ts2 link the ion pair IonP-N $\left[\mathrm{NH}_{4}^{+}\right]\left[\mathrm{BH}_{3} \mathrm{NH}_{2} \mathrm{BH}_{3}{ }^{-}\right]$with the $\mathrm{H}_{2}$ release products (Figure 12). The trans chain product DADB-p1 from the reaction involving IonP-N-ts1, whose geometry is shown in Figure 9, turns out to be the same as that from DADB. The shape of this TS is similar to that of DADB-ts1, in which a trans conformation results from a migration of the $\mathrm{NH}_{4}$ moiety. The $\mathrm{H}_{2}$ formation arises equally from a $\mathrm{B}-\mathrm{H} \cdots \mathrm{H}-\mathrm{N}$ interaction. The energy barrier associated with the process IonP-N $\rightarrow \mathbf{I o n P}$-N-ts1 is $17.7 \mathrm{kcal} / \mathrm{mol}, 2.8$ $\mathrm{kcal} / \mathrm{mol}$ smaller than that of DADB proceeding through DADBts1 (Figure 12 at the same level of theory). IonP-N-ts2 is too high an energy process to be considered.

Despite many attempts, we were not successful in identifying a direct reaction pathway linking the neutral dimer dim with either ion pair DADB or IonP-N, apart from those shown in Figure 8. There is a path to DADB that involves breaking the dimer interaction and starting from two monomers. The first step to form the two monomers requires $14.0 \mathrm{kcal} / \mathrm{mol}$. A second step is to break the $\mathrm{B}-\mathrm{N}$ bond, which requires an additional $25.9 \mathrm{kcal} / \mathrm{mol}$. An additional energy of only $1.9 \mathrm{kcal} / \mathrm{mol}$ is required to reach the value of $41.8 \mathrm{kcal} / \mathrm{mol}$ to reach TS $\mathbf{2 A B}$ ts3 leading to DADB. The actual reaction proceeds through the intermediate BBA-com $+\mathrm{NH}_{3}$, which is reached by a barrier of $12.4 \mathrm{kcal} / \mathrm{mol}$ starting from $2 \mathbf{A B}$. The energy barrier of 41.8 $\mathrm{kcal} / \mathrm{mol}$ is a few $\mathrm{kcal} / \mathrm{mol}$ smaller than the values of 44.5 and $50.0 \mathrm{kcal} / \mathrm{mol}$ for one-step elimination of one $\mathrm{H}_{2}$ (Figure 5) and two $\mathrm{H}_{2}$ molecules (Figure 5), respectively, from dim.

Thermal decomposition of $\mathbf{A B}$ has been experimentally shown to yield $\mathrm{H}_{2}$ at mild temperatures $\left(T \leq 80^{\circ} \mathrm{C}\right) .{ }^{5}$ Our results


Figure 12. Schematic energy profiles illustrating two different reaction pathways for $\mathrm{H}_{2}$ release from IonP-N. Relative energies in $\mathrm{kcal} / \mathrm{mol}$ were obtained from $\operatorname{CCSD}(\mathrm{T}) / \mathrm{aVTZ}+\mathrm{ZPE}$ calculations.
suggest that $\mathrm{H}_{2}$ release from ammonia borane under these conditions does not proceed through the hydrogen-bonded dimer dim. The dimer dim can generate one or two $\mathrm{H}_{2}$ molecules through direct pathways having energy barriers ranging from 44.5 to $50.0 \mathrm{kcal} / \mathrm{mol}$ with respect to the dimer (Figure 5). Instead, our results show that $\mathrm{H}_{2}$ can be released more readily if an ion pair is formed. Although DADB is $10.6 \mathrm{kcal} / \mathrm{mol}$ less stable than dim, the pathway for $\mathrm{H}_{2}$ elimination from DADB is characterized by a smaller energy barrier of $20.1 \mathrm{kcal} / \mathrm{mol}$. In the gas phase, the path that converts $\operatorname{dim}$ to DADB is has a barrier of $41.8 \mathrm{kcal} / \mathrm{mol}$ (through $\mathbf{2 A B}-\mathrm{ts} \mathbf{3}$ ). Such a conversion could be a lower-energy process in condensed media, ${ }^{15}$ so the barrier to lose $\mathrm{H}_{2}$ from the dimer would be reduced to 30.7 $\mathrm{kcal} / \mathrm{mol}$ in the absence of additional solid-state effects.

The second ion pair $\left[\mathrm{NH}_{4}{ }^{+}\right]\left[\mathrm{BH}_{3} \mathrm{NH}_{2} \mathrm{BH}_{3}{ }^{-}\right]$IonP- $\mathbf{N}$ is 16.4 $\mathrm{kcal} / \mathrm{mol}$ less stable than the neutral dimer and has a smaller energy barrier of $17.7 \mathrm{kcal} / \mathrm{mol}$ for $\mathrm{H}_{2}$ elimination. Again, the barrier for loss of $\mathrm{H}_{2}$ starting from dim via IonP-N would be $34.1 \mathrm{kcal} / \mathrm{mol}$, slightly higher than proceeding through DADB. These barriers are both below the unimolecular loss of $\mathrm{H}_{2}$ from $\mathrm{BH}_{3} \mathrm{NH}_{3} .{ }^{13}$ We found that the decomposition of both ion pairs produces the chain $\mathrm{BH}_{3} \mathrm{NH}_{2} \mathrm{BH}_{2} \mathrm{NH}_{3}$ structure, just as observed in the solid-state and ionic liquid studies. ${ }^{15,8 b}$

## Concluding Remarks

In the present study, we have investigated the reaction pathways and the molecular mechanisms for $\mathrm{H}_{2}$ release from
dimers of ammonia borane. In addition to the more conventional neutral dimer, we have found two ion pair isomers. One of them is DADB, which was experimentally identified following seeding of the neutral $(\mathbf{A B})_{2}$ dimer. The most stable gaseous phase dimer dim results from a head-to-tail cyclic interaction and is stabilized by $14.0 \mathrm{kcal} / \mathrm{mol}$ with respect to the two $\mathbf{A B}$ monomers. The neutral dimer dim could generate one $\mathrm{H}_{2}$ molecule through direct pathways having energy barriers near $45 \mathrm{kcal} / \mathrm{mol}$, substantially higher than for loss from the monomer. The pathway for loss of two $\mathrm{H}_{2}$ molecules is even higher near $50 \mathrm{kcal} / \mathrm{mol}$. Reactions of two monomers and of one monomer with separated $\mathrm{BH}_{3}$ and $\mathrm{NH}_{3}$ lead to either DADB or $\mathrm{H}_{2}$ formation with energy barriers of $\sim 26-28 \mathrm{kcal} / \mathrm{mol}$. Successive additions of two $\mathrm{NH}_{3}$ molecules to $\mathrm{B}_{2} \mathrm{H}_{6}$ also lead to DADB via an energetically lower-lying TS. Although DADB is $10.6 \mathrm{kcal} / \mathrm{mol}$ less stable than dim, the pathway for $\mathrm{H}_{2}$ elimination from DADB is characterized by a smaller energy barrier of $\sim 21 \mathrm{kcal} / \mathrm{mol}$. IonP-N is $16.4 \mathrm{kcal} / \mathrm{mol}$ less stable than the neutral dimer, and the energy barrier for loss of $\mathrm{H}_{2}$ is reduced to $\sim 18 \mathrm{kcal} / \mathrm{mol}$. The chain molecule $\mathrm{BH}_{3} \mathrm{NH}_{2} \mathrm{BH}_{2}-$ $\mathrm{NH}_{3}$, a key precursor for the formation of borane amine oligomers, was predicted to be generated when $\mathrm{H}_{2}$ is lost from either ion pair but not from the more stable dim. Our results are consistent with the experimental observation that seeding of $\mathrm{BH}_{3} \mathrm{NH}_{3}$ produces the diammoniate of diborane.

Acknowledgment. We thank Professor Michael McKee (Auburn University), Professor Rodger Nutt (Davidson College), and Dr. Donald Camaioni (Pacific Northwest National Laboratory) for valuable discussions. We thank the first two for suggesting the dissociative path to form DADB. Funding was provided in part by the Office of Energy Efficiency and Renewable Energy, Department of Energy, under the Hydrogen Storage Grand Challenge, Solicitation No. DE-PS36-03GO93013. This work was done as part of the Chemical Hydrogen Storage Center. D.A.D. is indebted to the Robert Ramsay Endowment of the University of Alabama. V.S.N. thanks the Belgian Technical Cooperation Agency for a doctoral scholarship. M.T.N. is grateful to the FWO-Vlaanderen for supporting his sabbatical leave at the University of Alabama.

Supporting Information Available: $\operatorname{CCSD}(\mathrm{T})$ total energies, MP2 ZPEs, Cartesian coordinates of the MP2/aVTZoptimized geometries, and atomization energy components for the ion pairs and the ions $\mathrm{BH}_{3} \mathrm{NH}_{2} \mathrm{BH}_{3}{ }^{-}$and $\mathrm{NH}_{3} \mathrm{BH}_{2} \mathrm{NH}_{3}{ }^{+}$and figures showing important MP2/aVTZ geometry parameters of the transition structures dim-ts6, dim-ts7, and dim-ts8 for the release of one $\mathrm{H}_{2}$ molecule from $\operatorname{dim}$ and the various products. This material is available free of charge via the Internet at http:// pubs.acs.org.

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